

$$\delta_r v(x) = v(x+r) - v(x)$$

$$\delta_r v(x) = v_0 \left(\frac{r}{L_0} \right)^{h(x)}$$

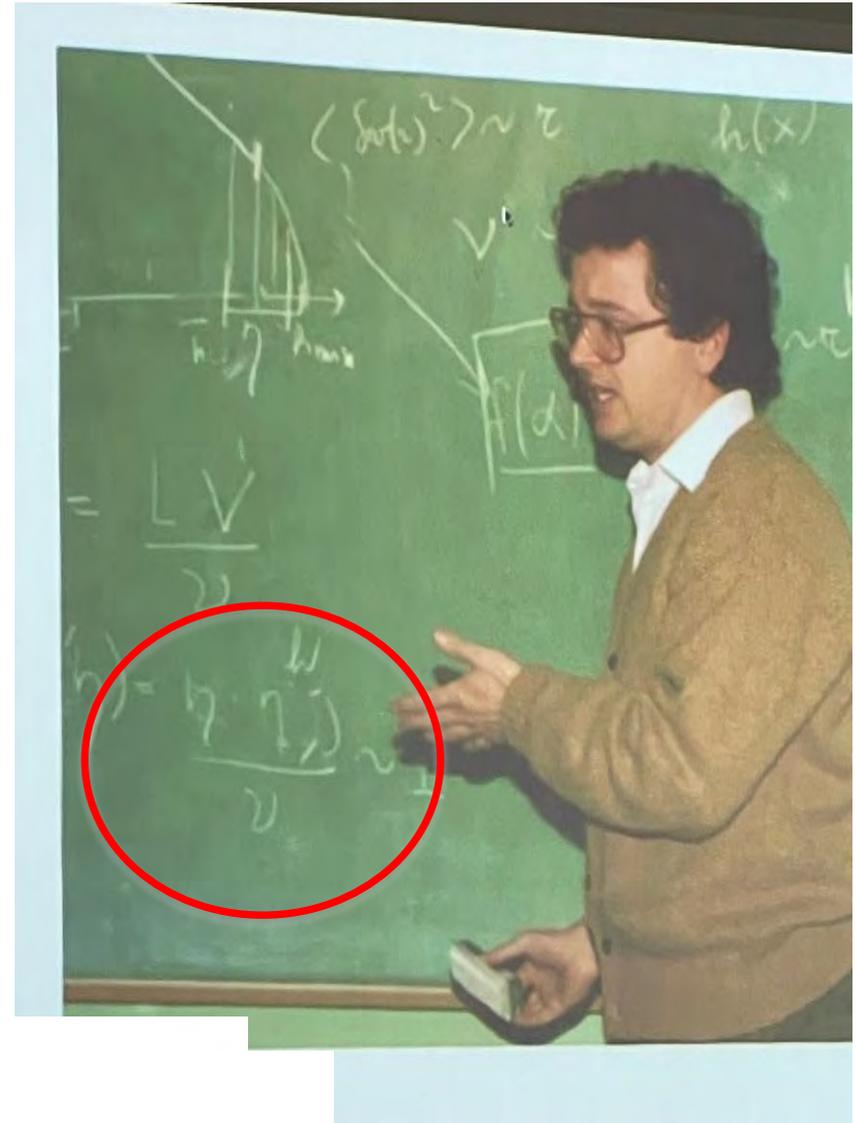
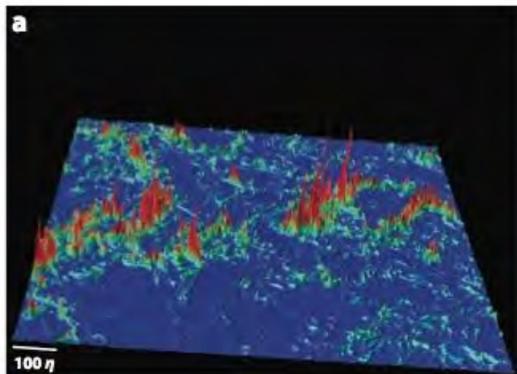
$$\frac{\delta_\eta v \eta}{\nu} \sim O(1)$$

$$\frac{v_0 (\eta/L_0)^h \eta}{\nu} \sim O(1)$$

$$\eta/L_0 \sim \nu^{\frac{1}{1+h(x)}} \sim Re^{-\frac{1}{1+h(x)}}$$

-INFINITELY MANY ANOMALOUS EXPONENTS

-INFINITELY MANY UV CUTOFFS



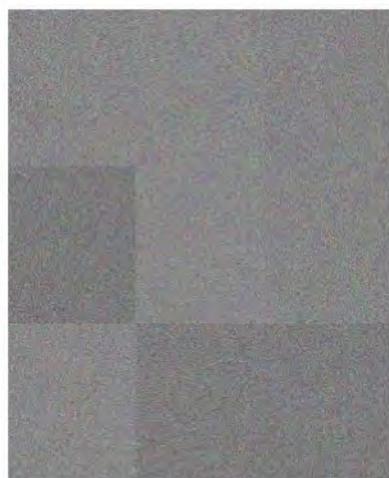
Degrees of freedom of turbulence

Giovanni Paladin and Angelo Vulpiani

Phys. Rev. A **35**, 1971(R) – Published 1 February 1987



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HOLLYWOOD STARS



LES HOUCHE-2024

Data driven tools for Lagrangian Turbulence

CREDITS: T. LI, F. BONACCORSO, M. BUZZICOTTI, M. SCARPOLINI



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Ministero dell'Università e della Ricerca



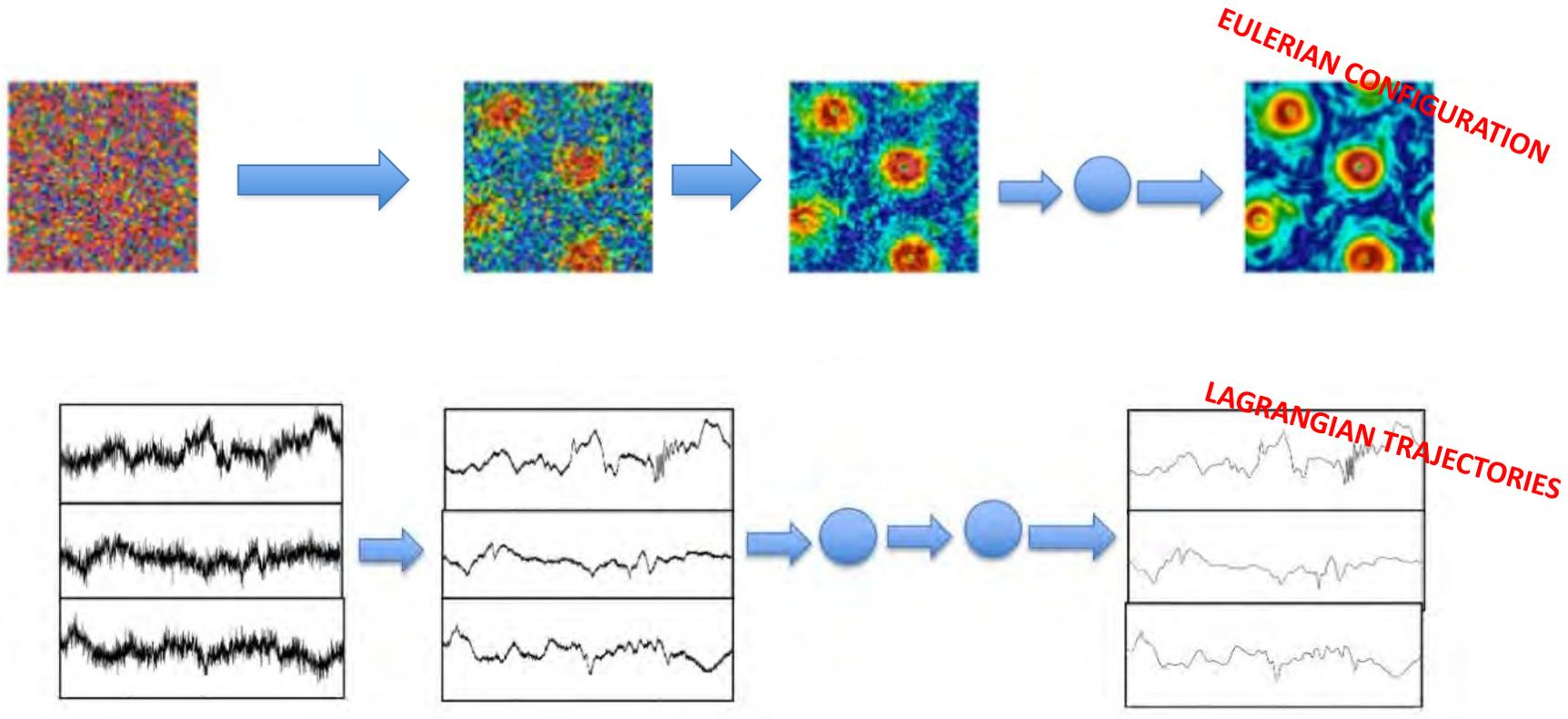
Italiadomani
PIANO NAZIONALE DI RIPRESA E RESILIENZA



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LES HOUCHES-2024

Data driven tools for Lagrangian Turbulence

CREDITS: T. LI, F. BONACCORSO, M. BUZZICOTTI, M. SCARPOLINI

STOCHASTIC MODELS FOR LAGRANGIAN TURBULENCE: WHY?

T. Li, LB, F. Bonaccorso, M. Scarpolini, M. Buzzicotti.

Synthetic Lagrangian Turbulence by Generative Diffusion Models. [arXiv:2307.08529](https://arxiv.org/abs/2307.08529) – in press Nature Machine Intelligence (2024)

GENERATION OF LARGE SYNTHETIC DATA-BASE FOR

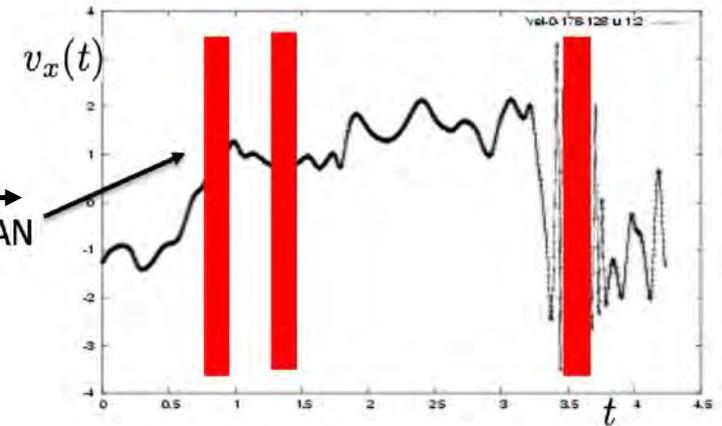
(I) RANKING OF PHYSICS FEATURES

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- (A) SINGLE, TWO AND MULTI-PARTICLES DISPERSION
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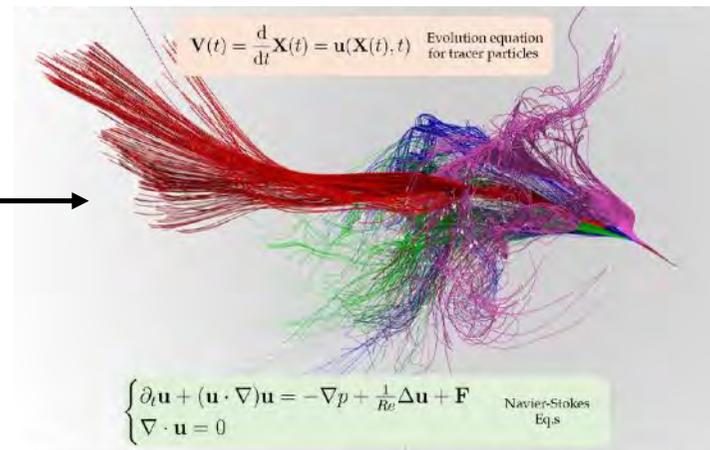
(III) DATA ASSIMILATION/IMPUTATION FROM MISSING FIELD/EXPERIMENTAL OBERVATION

LAGRANGIAN

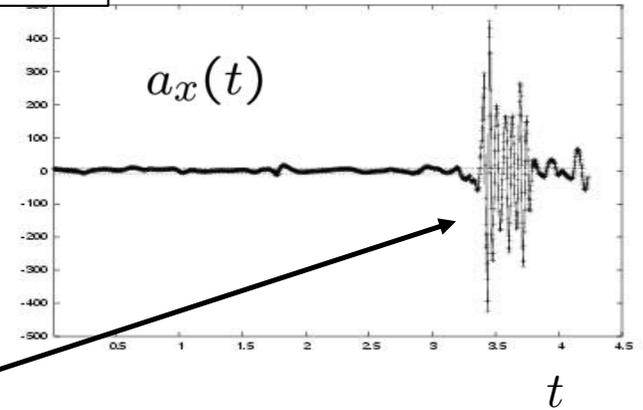
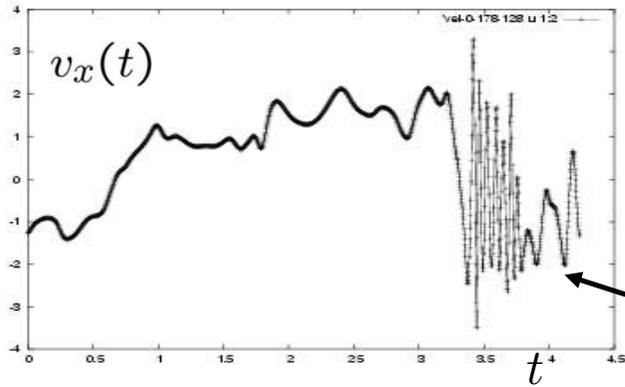


(IV) CLASSIFICATION/INFERRAL OF MISSING/INTERNAL PROPERTIES:

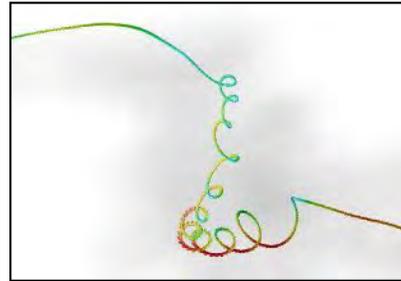
- i. INERTIA
- ii. SHAPE
- iii. ACTIVE DEGREES OF FREEDOM
- iv.



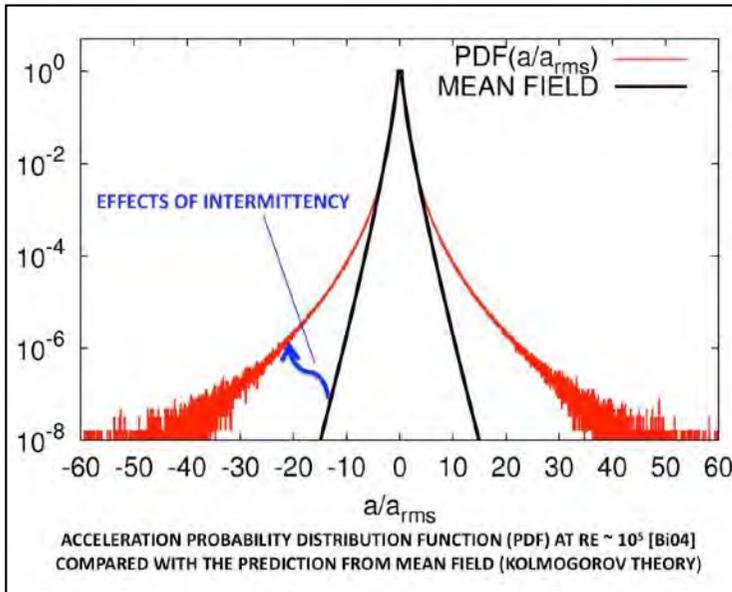
$$\begin{cases} \mathbf{a} = \partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p + \nu \Delta \mathbf{u} + \mathbf{f} \\ \nabla \cdot \mathbf{u} = 0 \end{cases}$$



EXTREME EVENTS



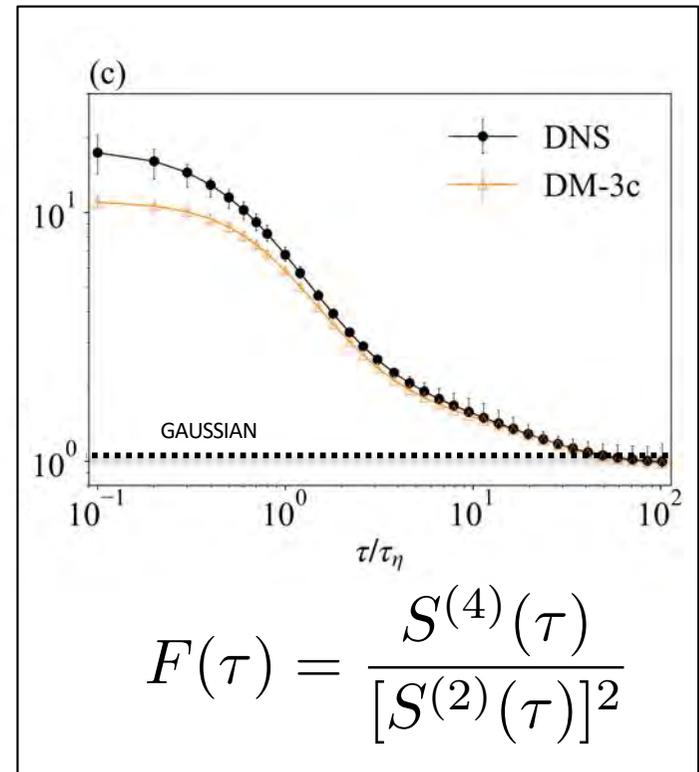
$$S_i^{(p)}(\tau) = \langle [v_i(t + \tau) - v_i(t)]^p \rangle$$



La Porta, G.A. Voth, A.M. Crawford, J. Alexander et al. Fluid particle accelerations in fully developed turbulence. *Nature*, 409(6823), 1017 (2001)

N. Mordant, P. Metz, O. Michel and J.F. Pinton. Measurement of Lagrangian velocity in fully developed turbulence. *Phys. Rev. Lett.* 87(21), 214501 (2001)

F. Toschi and E. Bodenschatz. Lagrangian Properties of Particles in Turbulence. *Annu. Rev. Fluid Mech.* 41, 375 (2009)



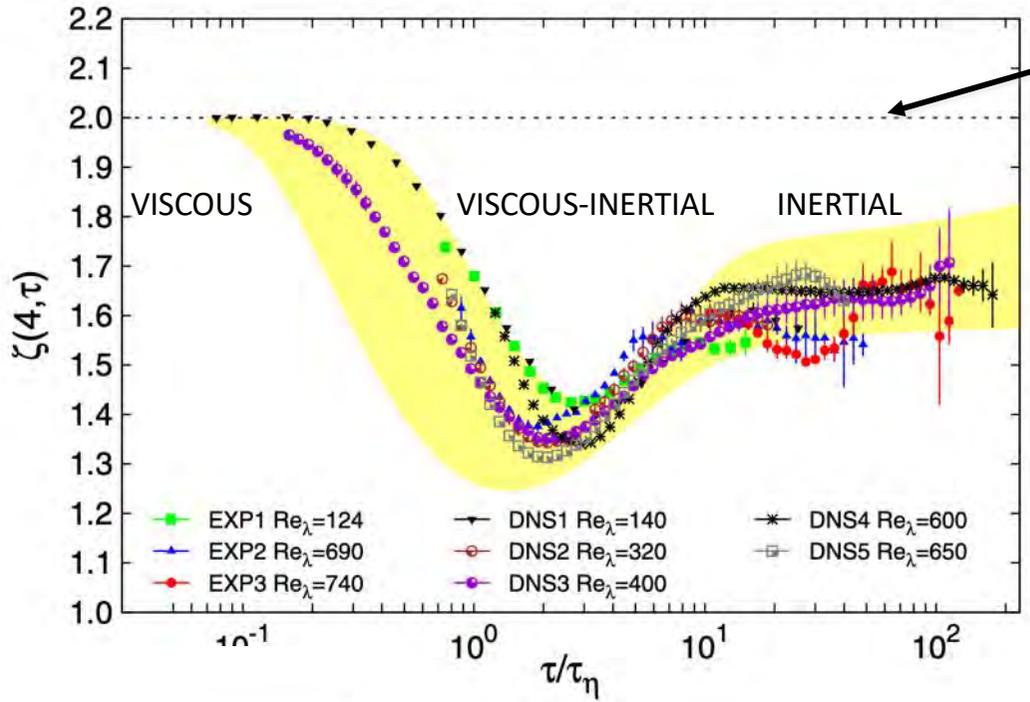
$$S_i^{(p)}(\tau) = \langle [v_i(t + \tau) - v_i(t)]^p \rangle$$

$$\zeta(4, \tau) = \frac{d \log S^{(4)}(\tau)}{d \log S^{(2)}(\tau)}$$

Universal Intermittent Properties of Particle Trajectories in Highly Turbulent Flows

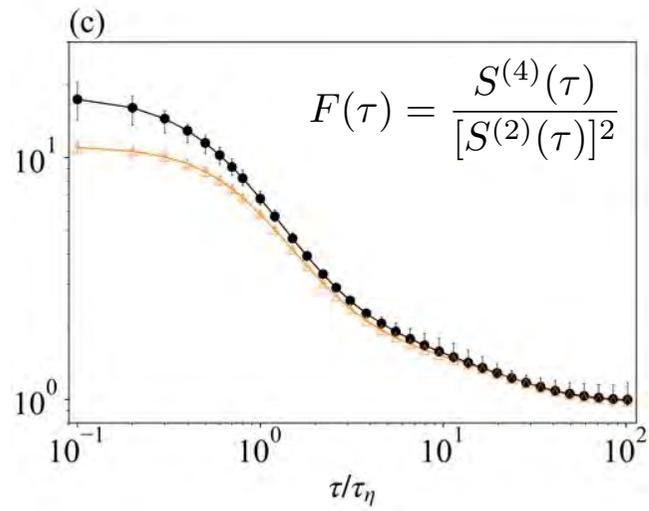
A. Arnéodo,¹ R. Benzi,² J. Berg,³ L. Biferale,^{4,*} E. Bodenschatz,⁵ A. Busse,⁶ E. Calzavarini,⁷ B. Castaing,¹ M. Cencini,^{8,*} L. Chevillard,¹ R. T. Fisher,⁹ R. Grauer,¹⁰ H. Homann,¹⁰ D. Lamb,⁹ A. S. Lanotte,^{11,*} E. Lévêque,¹ B. Lüthi,¹² J. Mann,³ N. Mordant,¹³ W.-C. Müller,⁶ S. Ott,³ N. T. Ouellette,¹⁴ J.-F. Pinton,¹ S. B. Pope,¹⁵ S. G. Roux,¹ F. Toschi,^{16,17,*} H. Xu,⁵ and P. K. Yeung¹⁸

(International Collaboration for Turbulence Research)



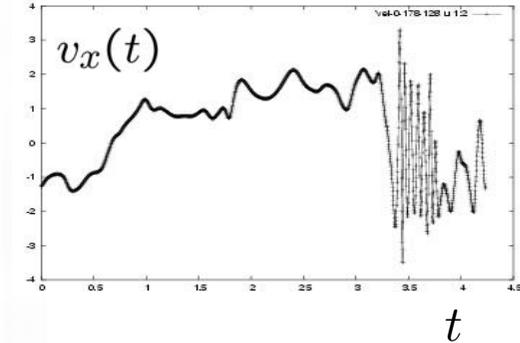
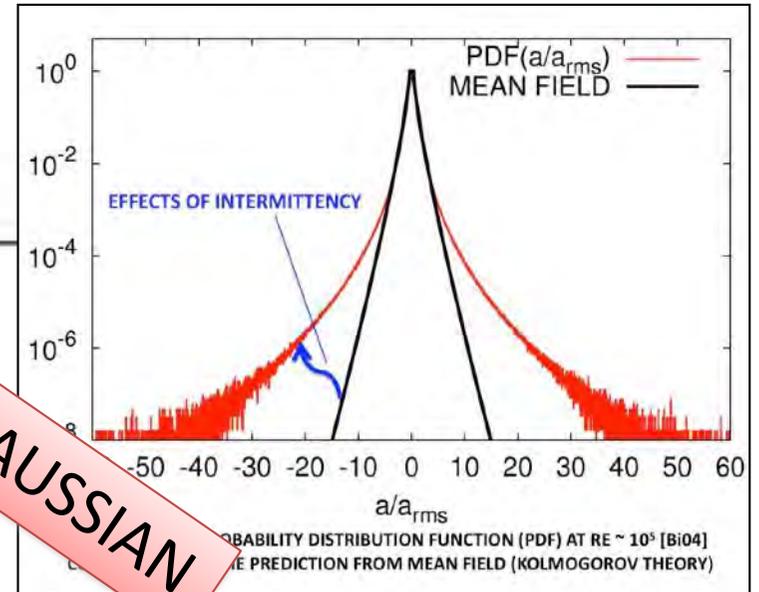
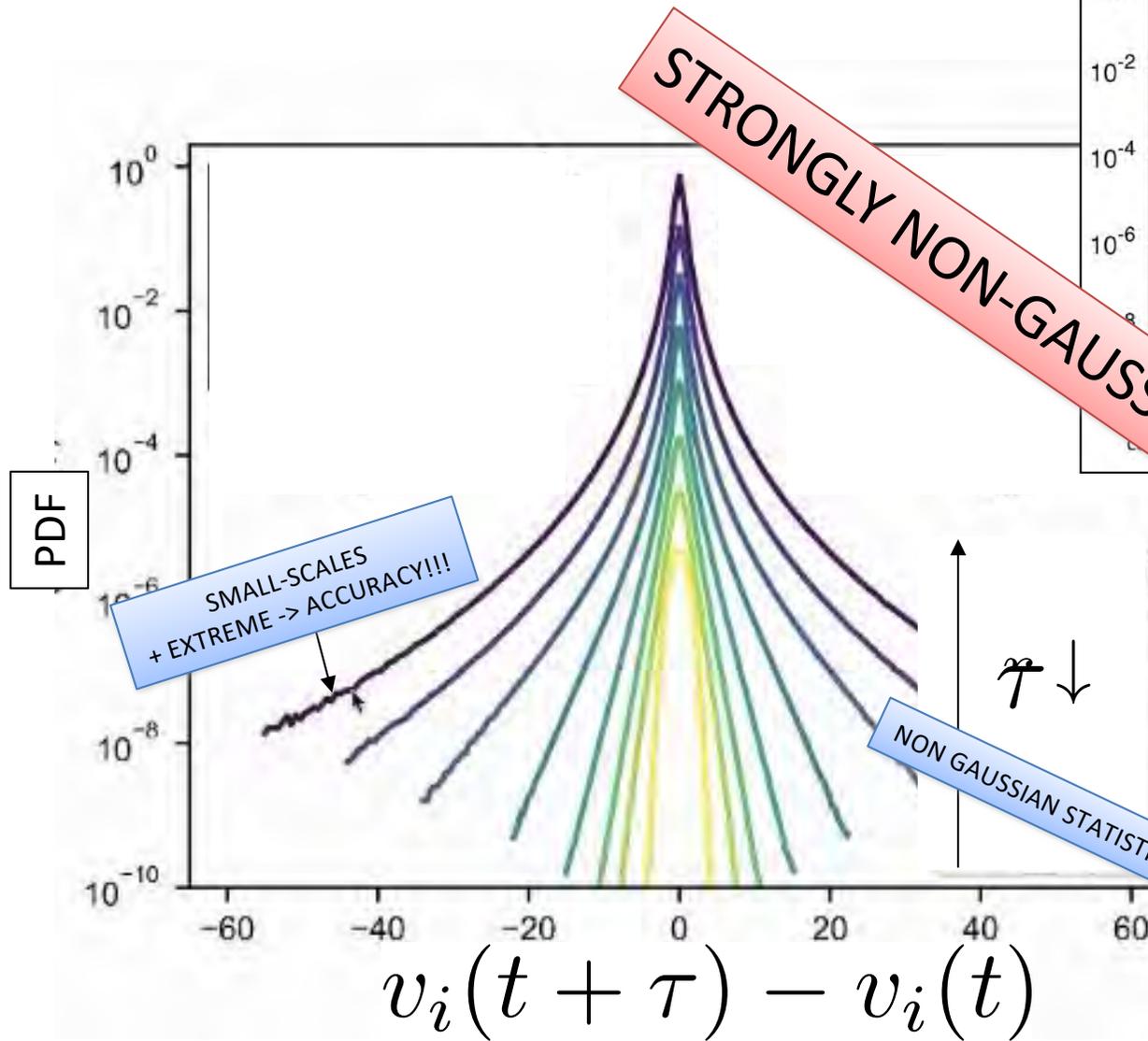
GAUSSIAN – NO INTERMITTENCY

- UNIVERSALITY
- SCALE-BY-SCALE INTERMITTENCY
- VISCOUS-SCALE FLUCTUATIONS
- MF-PREDICTION



M. Borgas "The multifractal Lagrangian nature of turbulence", PTRSA 342, 379 (1993)
 G. K. Batchelor. "Pressure fluctuations in isotropic turbulence" Proc. Camb. Philos. Soc. 47, 359 (1951)
 G. Paladin and A. Vulpiani, "Degrees of freedom of turbulence," Phys. Rev. A 35, 1971 (1987)
 C. Meneveau, "Transition between viscous and inertial-range scaling of turbulence structure functions" Phys. Rev. E 54, 3657 (1996)

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Chevillard, L., Garban, C., Rhodes, R. & Vargas, V. On a skewed **and multifractal unidimensional random field**, as a probabilistic representation of kolmogorov's views on turbulence. In Annales Henri Poincaré, vol. 20, 3693–3741 (Springer, 2019).

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Zamansky, R. **Acceleration scaling and stochastic dynamics** of a fluid particle in turbulence. Physical Review Fluids 7, 084608 (2022).

Lubcke, J, Friedrich, J., Grauer, R. **Stochastic interpolation** of sparsely sampled time series by a superstatistical random process and its synthesis in Fourier and wavelet space. J. Phys. Complex. 4 015005 (2023)

Diffusion Models

Training set: a set of images $\vec{a}^\mu \in \mathbb{R}^N$ $\mu = 1, \dots, P$
N is the dimension of the data, P their number

Langevin equation for an Ornstein-Uhlenbeck process

$$\frac{d\vec{x}}{dt} = -\vec{x} + \vec{\eta}(t) \quad \langle \eta_i(t) \eta_j(t') \rangle = 2T \delta_{ij} \delta(t - t')$$

$\vec{x}^\mu(t = 0) = \vec{a}^\mu$ It transforms the data in iid Gaussian $\mathcal{N}(0, 1)$ at $t \gg 1$

$$P_t(\vec{x}) = \int d\vec{a} P_0(\vec{a}) \frac{1}{\sqrt{2\pi\Delta_t}^N} \exp\left(-\frac{1}{2} \frac{(\vec{x} - \vec{a}e^{-t})^2}{\Delta_t}\right) = \int d\vec{a} P_t(\vec{a}, \vec{x})$$

$\Delta_t = T(1 - e^{-2t})$



Score function provides the force field to go back in time

$$\mathcal{F}_i(\vec{x}, t) = \frac{\partial \log P_t(\vec{x})}{\partial x_i} \quad -\frac{dy_i}{dt} = y_i + 2T \mathcal{F}_i(y, t) + \eta_i(t)$$

DIFFUSION MODELS

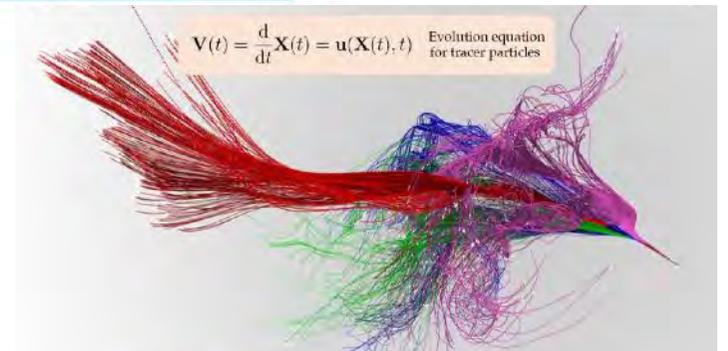
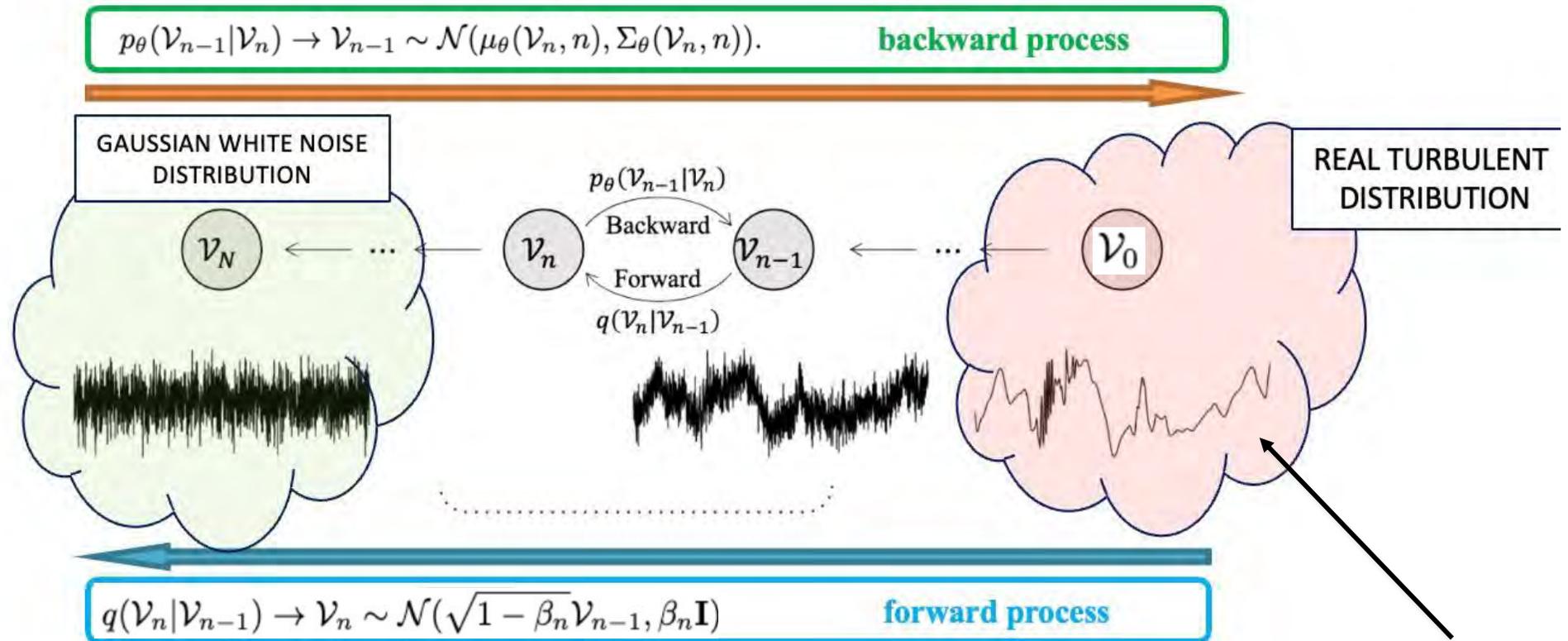
'Synthetic Lagrangian Turbulence: all you need is Diffusion Models'

T. Li, L.B, F. Bonaccorso, M. Scarpolini and M. Bucciotti (arXiv:2307.08529 2024, Nature Machine Intelligence in press)

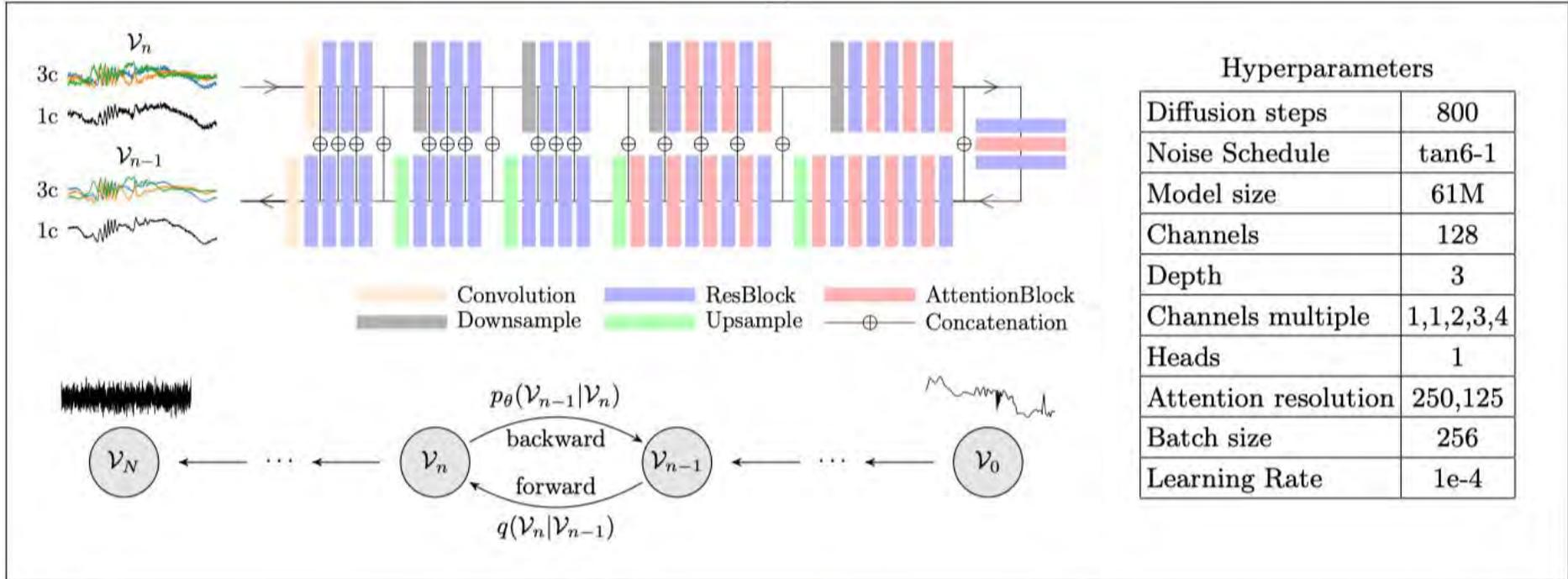
[Sohl-Dickstein et al., Deep Unsupervised Learning using Nonequilibrium Thermodynamics, ICML 2015](#)

[Ho et al., Denoising Diffusion Probabilistic Models, NeurIPS 2020](#)

[Song et al., Score-Based Generative Modeling through Stochastic Differential Equations, ICLR 2021](#)



(a)

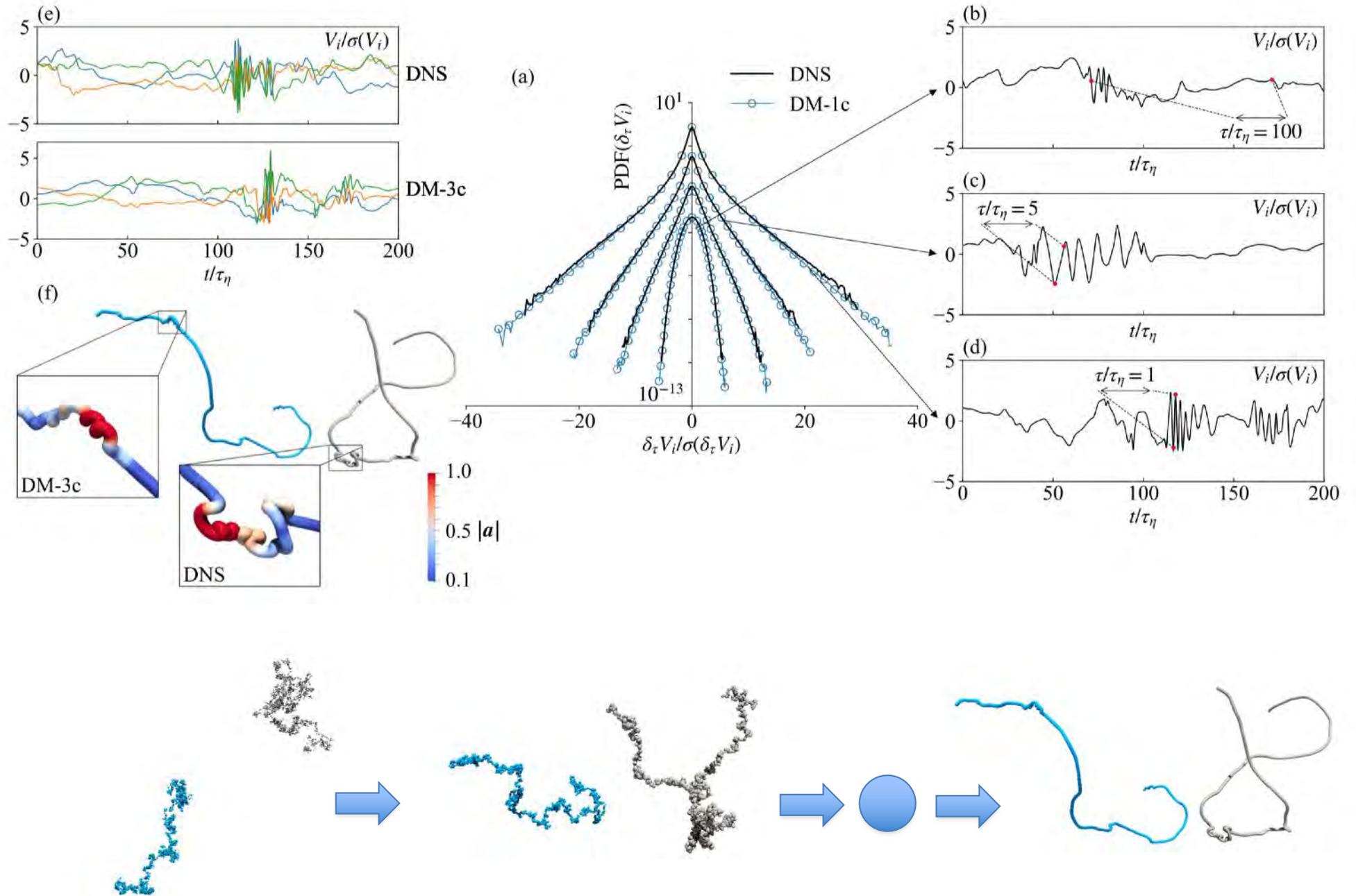


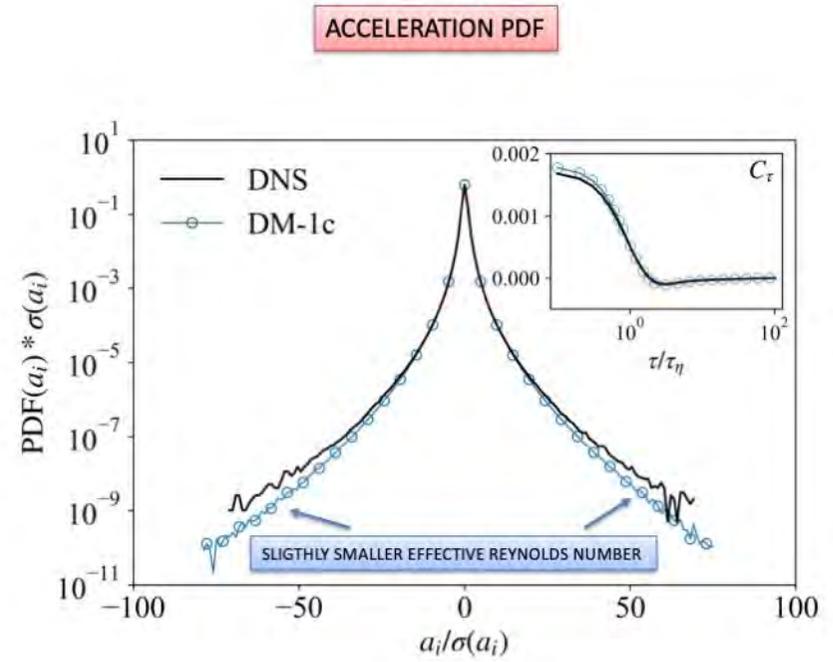
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[Ho et al., Denoising Diffusion Probabilistic Models, NeurIPS 2020](#)

[Song et al., Score-Based Generative Modeling through Stochastic Differential Equations, ICLR 2021](#)

$$\delta_\tau V_i(t) = V_i(t + \tau) - V_i(t),$$



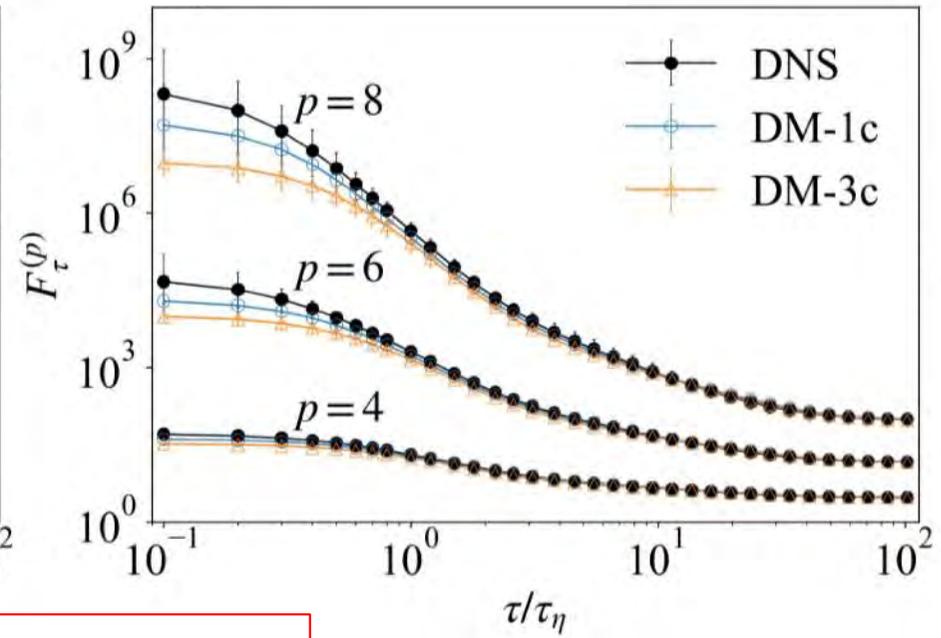
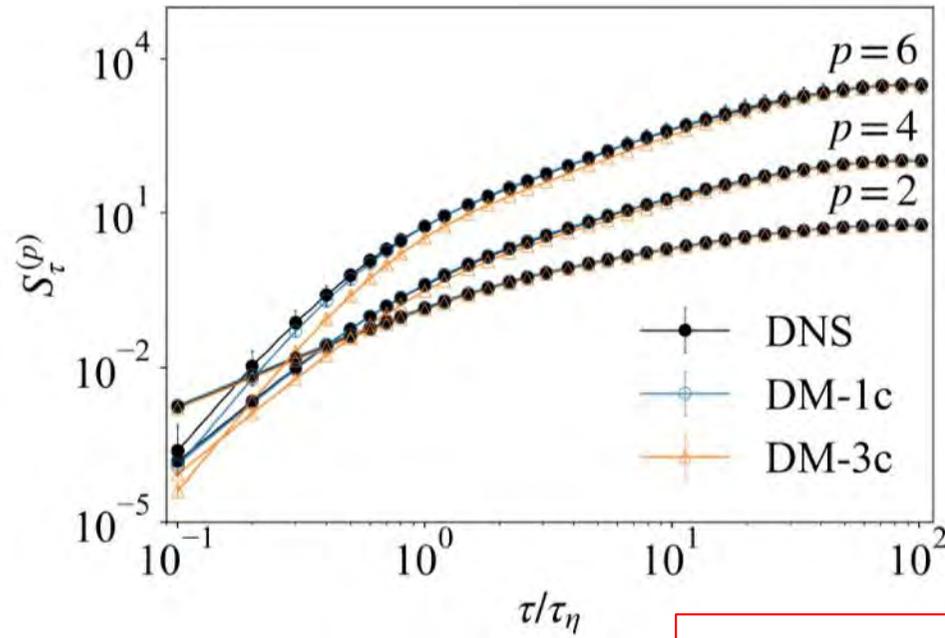


LAGRANGIAN STRUCTURE FUNCTIONS

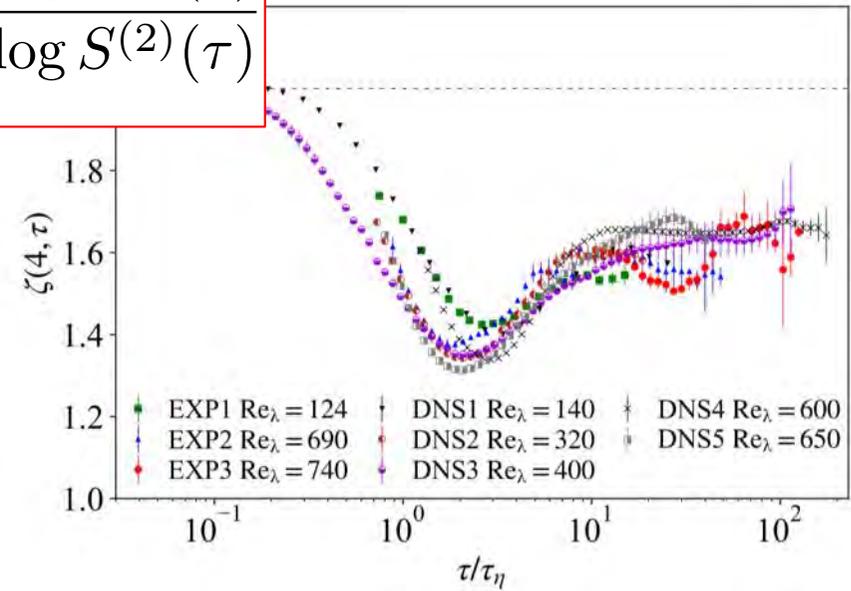
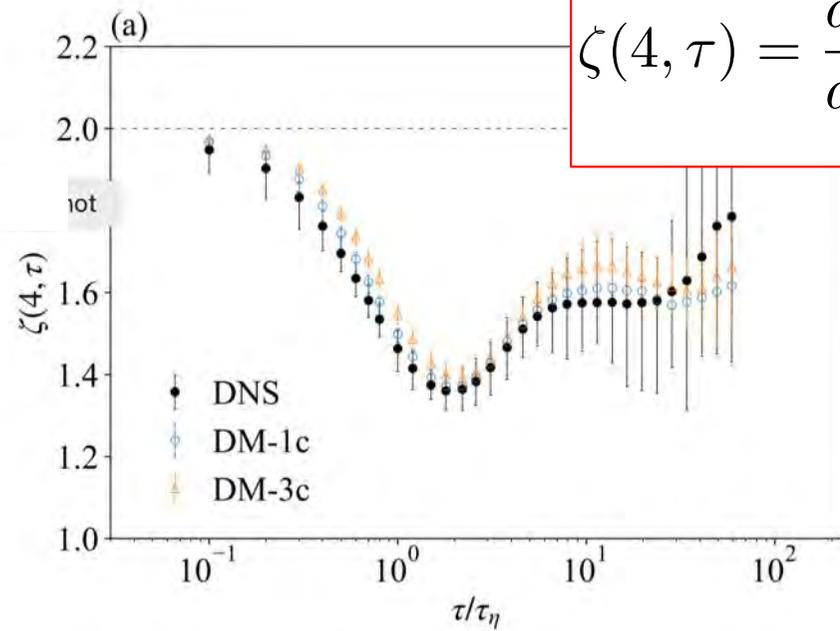
GENERALIZED FLATNESS

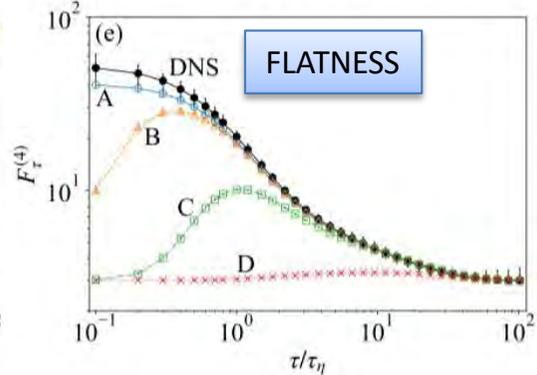
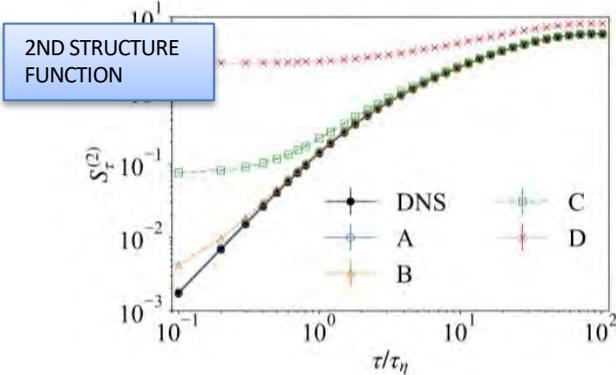
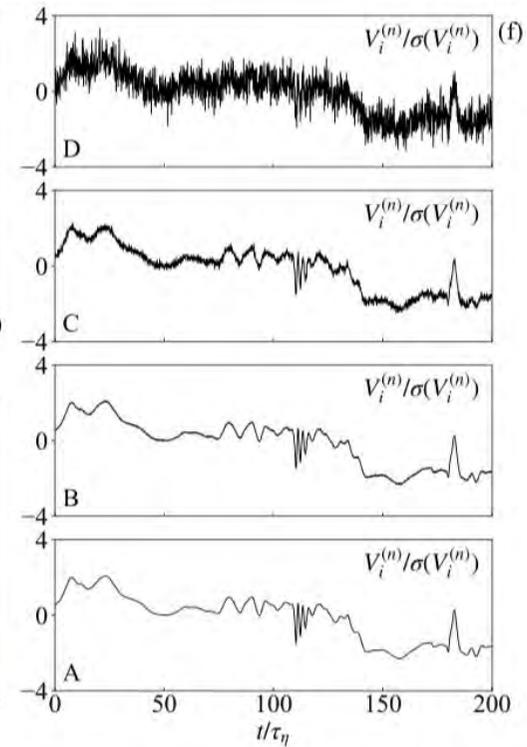
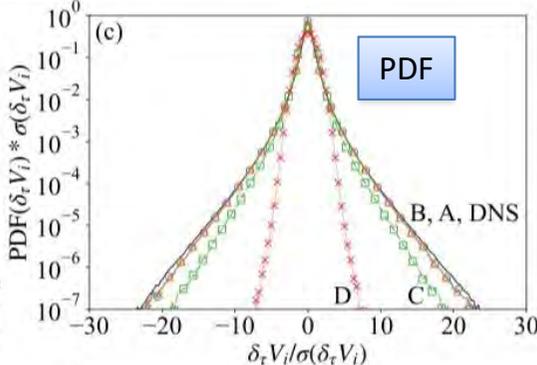
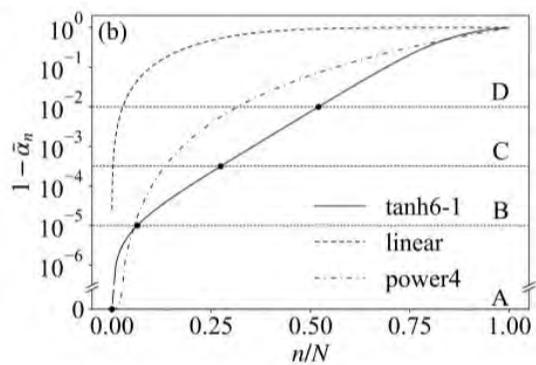
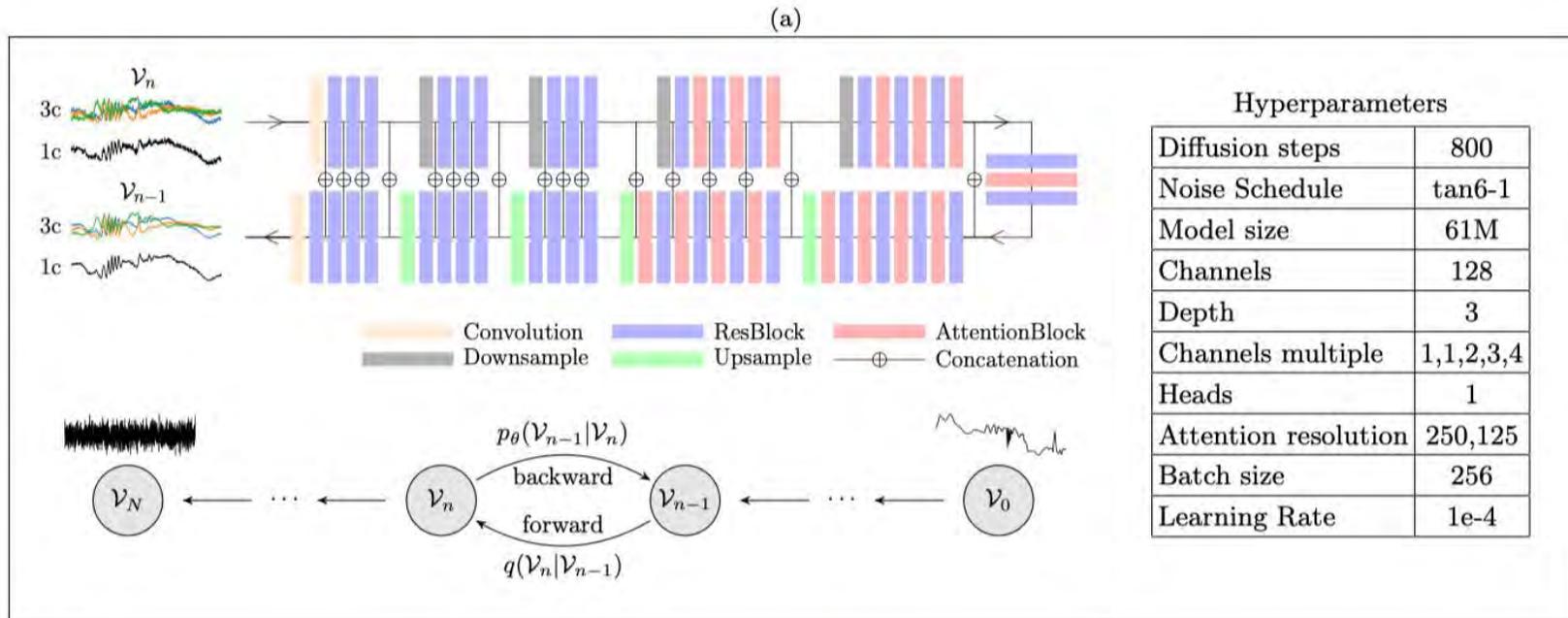
$$S_\tau^{(p)} = \langle (\delta_\tau V_i)^p \rangle$$

$$F_\tau^{(p)} = S_\tau^{(p)} / [S_\tau^{(2)}]^{p/2}$$



$$\zeta(4, \tau) = \frac{d \log S^{(4)}(\tau)}{d \log S^{(2)}(\tau)}$$





STOCHASTIC MODELS FOR LAGRANGIAN TURBULENCE: WHY?

T. Li, LB, F. Bonaccorso, M. Scarpolini, M. Buzzicotti.

Synthetic Lagrangian Turbulence by Generative Diffusion Models. [arXiv:2307.08529](https://arxiv.org/abs/2307.08529) – in press Nature Machine Intelligence (2024)

GENERATION OF LARGE SYNTHETIC DATA-BASE FOR

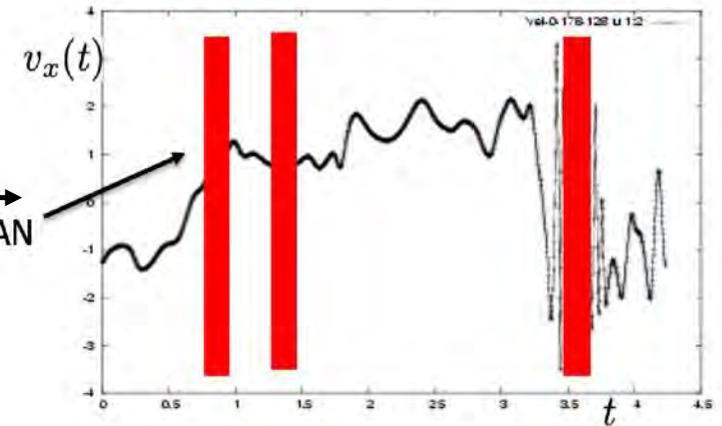
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- (D) MULTI-TIME MULTI-SCALE TURBULENT FLUCTUATIONS

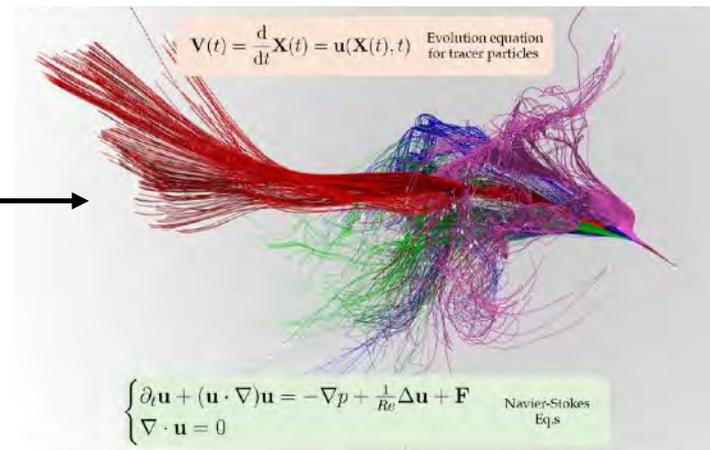
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LAGRANGIAN



(IV) CLASSIFICATION/INFERRAL OF MISSING/INTERNAL PROPERTIES:

- i. INERTIA
- ii. SHAPE
- iii. ACTIVE DEGREES OF FREEDOM
- iv.



IMPUTATION, CONDITIONAL TEXT GENERATION

Since Kolmogorov's conjecture, in this so-called range have been expected to follow universal for which theoretical predictions

1. Since Kolmogorov's conjecture, **scaling laws** in this so-called **inertial range** have been expected to follow universal **behavior** for which theoretical predictions **exist**.

2. Since Kolmogorov's conjecture, **turbulent flows** in this so-called **inertial range** have been expected to follow universal **scaling laws** for which theoretical predictions **abound**.

GT: Since Kolmogorov's conjecture, the **velocity difference statistics** in this so-called **inertial range** have been expected to follow universal **power laws** for which theoretical predictions **have been refined over the years**

UNCONDITIONED TEXT GENERATION

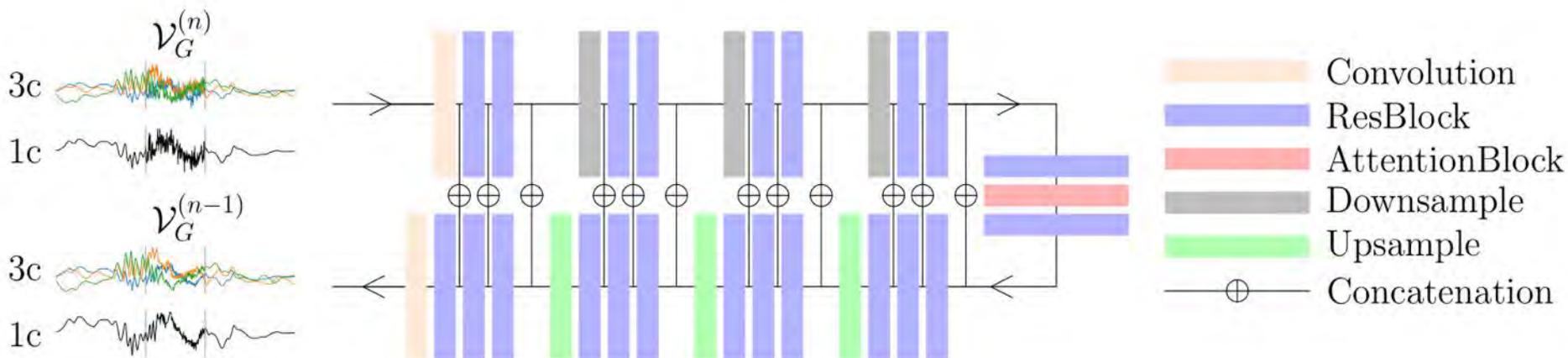
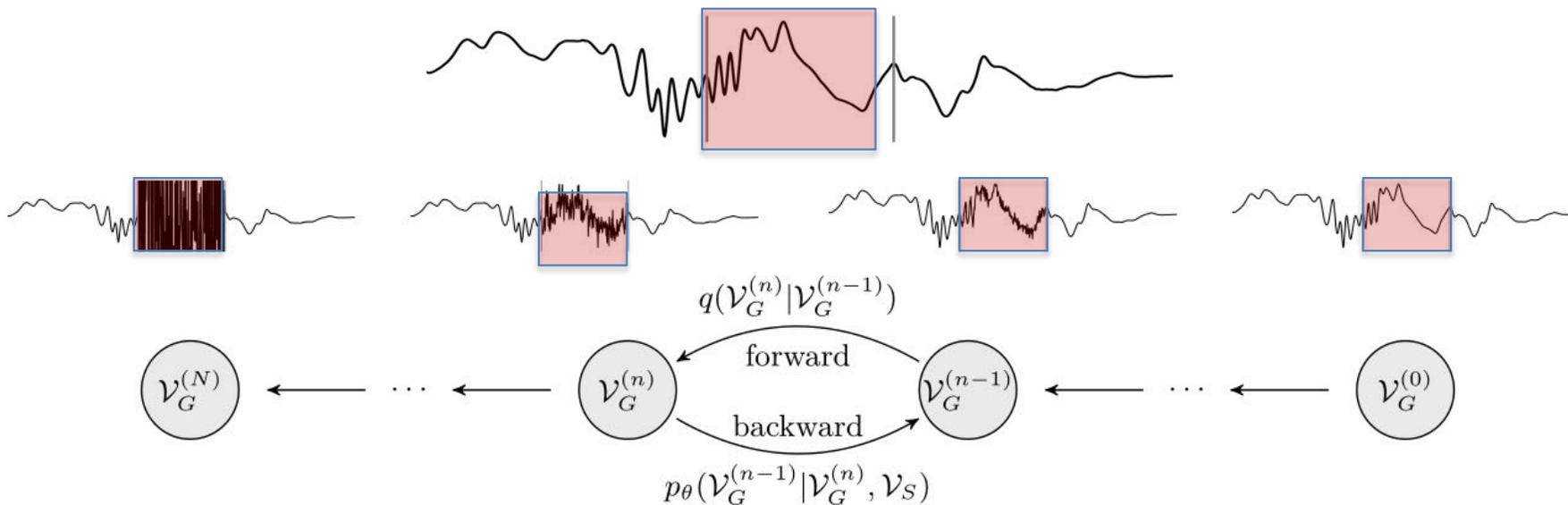
Q: TELL MORE ABOUT UNIVERSALITY OF SMALL-SCALE FLUCTUATIONS IN TURBULENCE

Small-scale fluctuations in turbulence **are not entirely universal across all scales and flow conditions**. While certain statistical properties, such as the energy cascade from large to small scales, exhibit universal behavior as predicted by **theories** like **Kolmogorov's**, the specifics of small-scale fluctuations can vary depending on factors such as **the Reynolds number, flow geometry, and presence of external forces**.

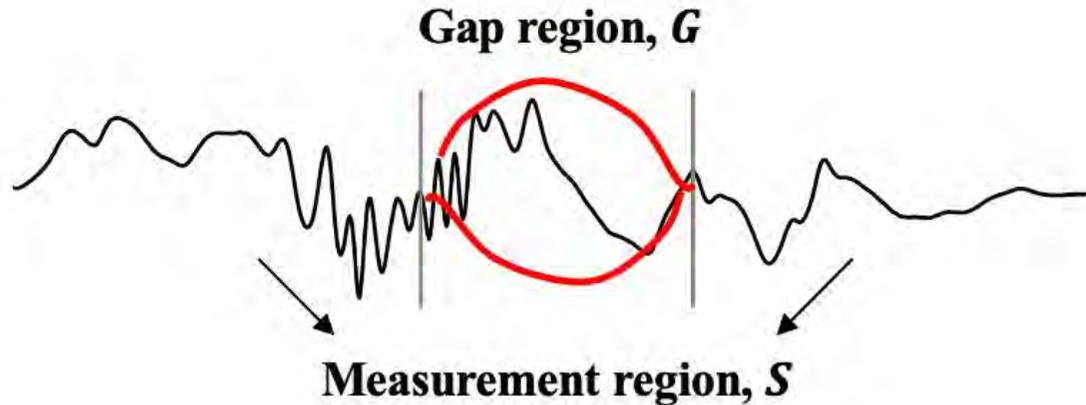
However, there are certain aspects of small-scale turbulence that can be considered **more universal**. For example, in the **inertial subrange of high Reynolds number turbulence, the statistical behavior of small-scale fluctuations tends to follow power-law scaling consistent with Kolmogorov's theories**. Yet, the precise values of parameters such as the **Kolmogorov constant or the intermittency corrections can vary in different flows**. Therefore, while there are universal features in small-scale turbulence, there are also non-universal aspects that depend on the specific conditions of the flow.

IMPUNTATION OF LAGRANGIAN TRAJECTORIES: CONDITIONAL DM

Gap size: $T_G/\tau_\eta = 50, T_G/T_I = 1/4$



Gaussian Process Regression (GPR)



Training process:

$$\mathcal{V}_S = \{V(t_{s_1}), V(t_{s_2}), \dots, V(t_{s_{N(S)}}) \mid t_{s_i} \in S\}$$

$$\mathcal{V}_G = \{V(t_{g_1}), V(t_{g_2}), \dots, V(t_{g_{N(G)}}) \mid t_{g_i} \in G\}$$

$$\begin{bmatrix} \mathcal{V}_S \\ \mathcal{V}_G \end{bmatrix} \sim \mathcal{N}\left(\mathbf{0}, \begin{bmatrix} C_{SS} & C_{SG} \\ C_{GS} & C_{GG} \end{bmatrix}\right)$$

$$(C_{TT})_{ij} = \langle V(t_i)V(t_j) \rangle$$

computing the covariance with training data

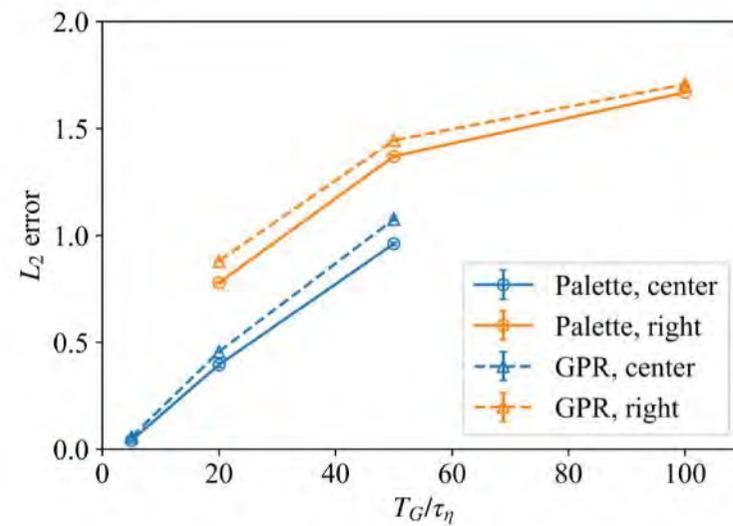
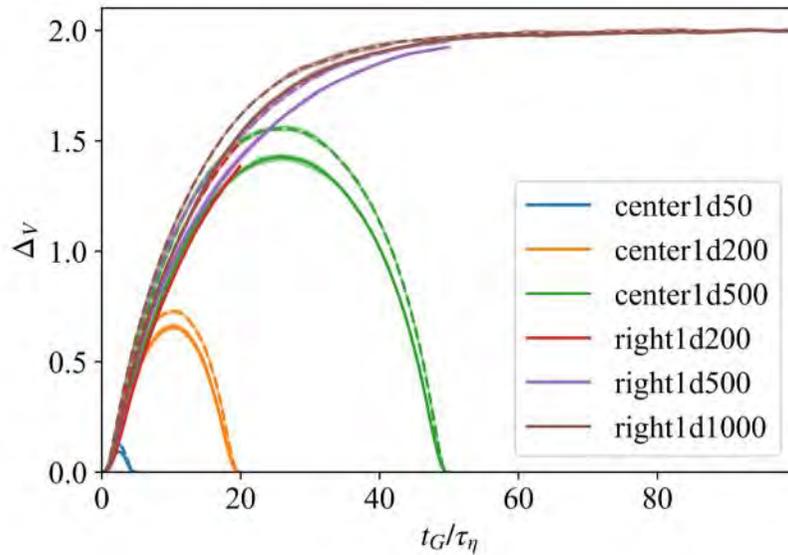
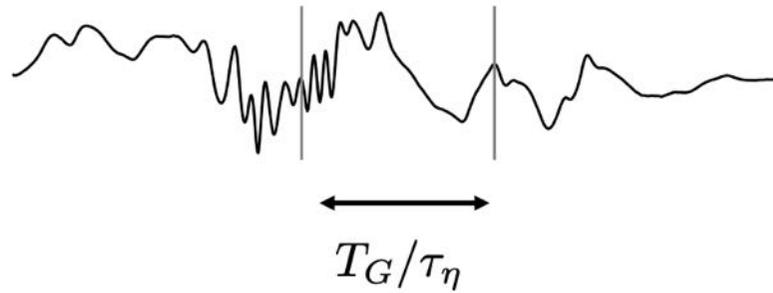
Reconstruction with measurement, \mathcal{V}_S :

$$\mathcal{V}_G \mid \mathcal{V}_S \sim \mathcal{N}(\mu_G, \Sigma_{GG})$$

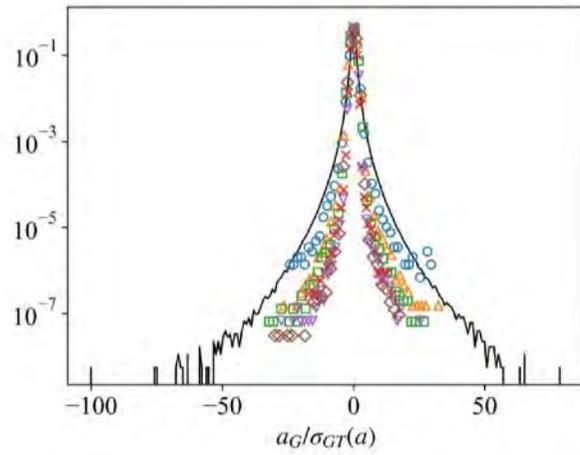
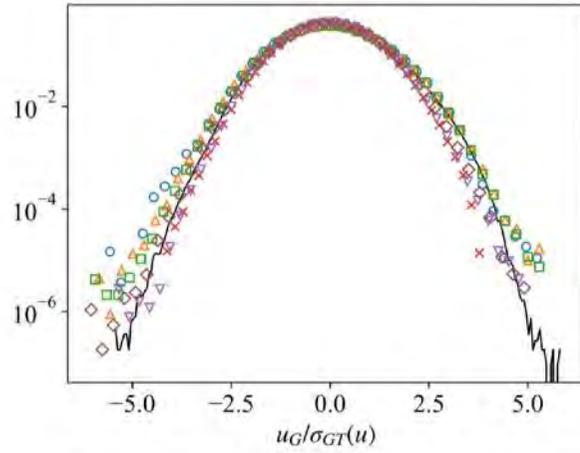
$$\mu_G = C_{GS}C_{SS}^{-1}\mathcal{V}_S$$

$$\Sigma_{GG} = C_{GG} - C_{GS}C_{SS}^{-1}C_{SG}$$

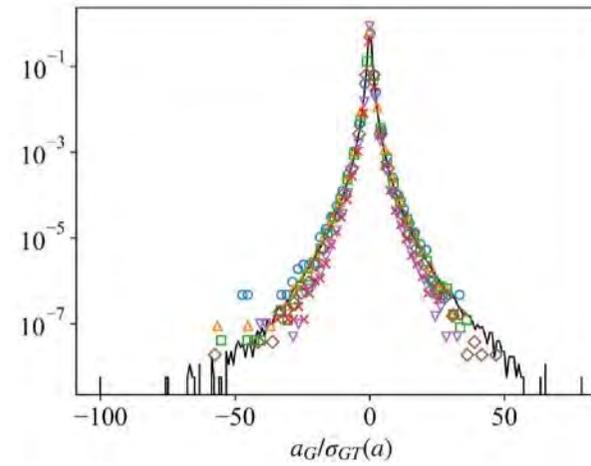
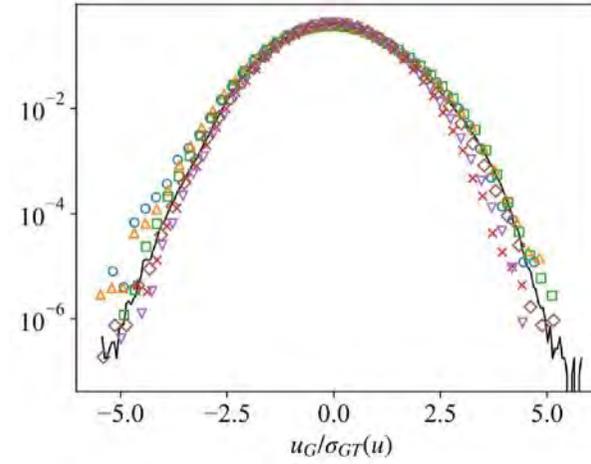
$$\Delta V(t_G) = \frac{\langle [V_p(t_G) - V_{gt}(t_G)]^2 \rangle}{\sqrt{\langle [V_p]^2 \rangle \langle [V_{gt}]^2 \rangle}}$$



GAUSSIAN PROCESS REGRESSION



DIFFUSION MODEL



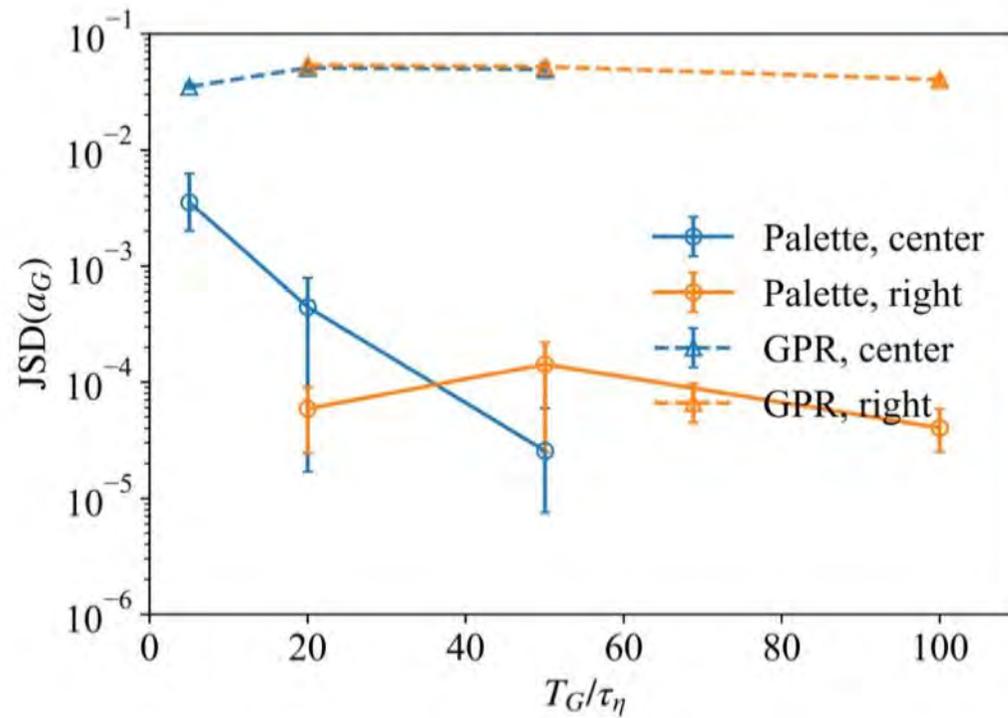
KULLBACK-LEIBLER

$$D(P \parallel Q) = \int_{-\infty}^{\infty} P(x) \log \left(\frac{P(x)}{Q(x)} \right) dx$$

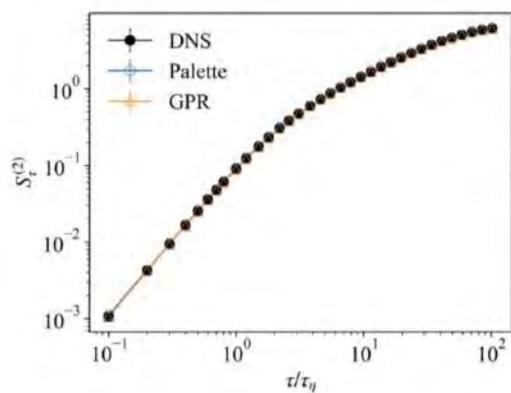
$$\xrightarrow{M = \frac{1}{2}(P+Q)}$$

JENSEN-SHANNON

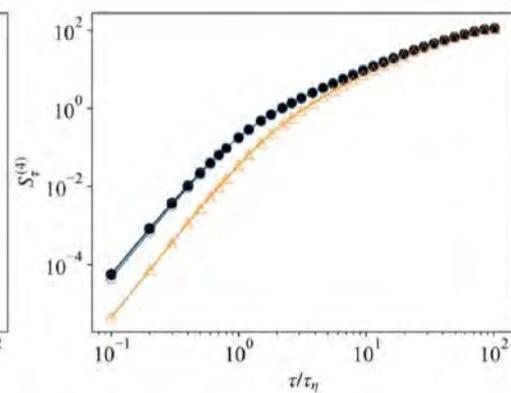
$$JSD(P \parallel Q) = \frac{1}{2}D(P \parallel M) + \frac{1}{2}D(Q \parallel M),$$



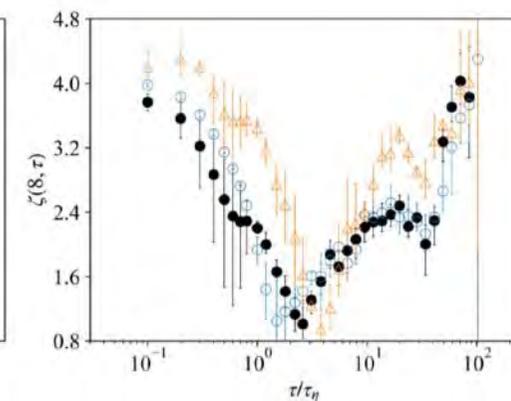
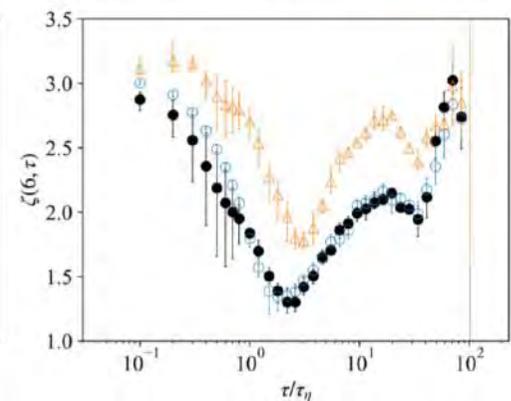
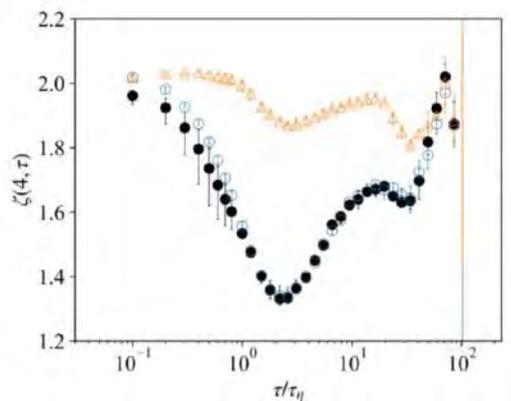
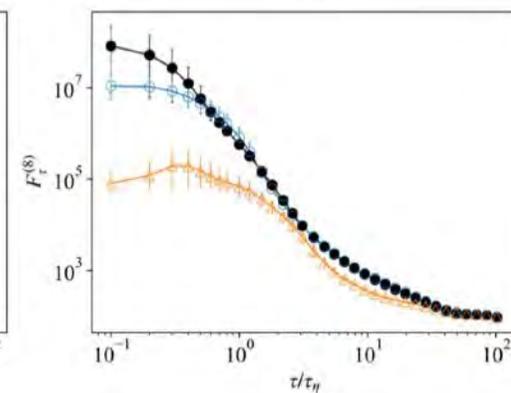
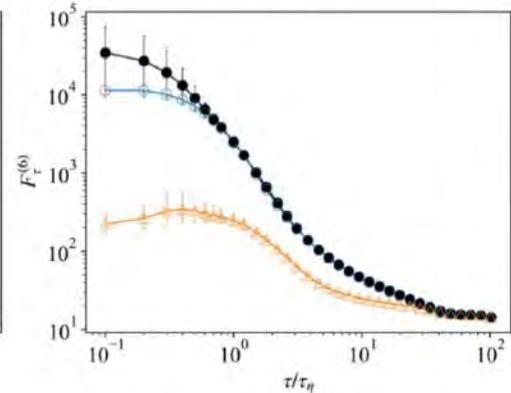
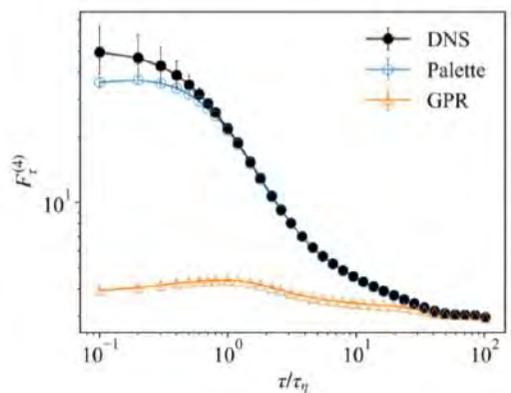
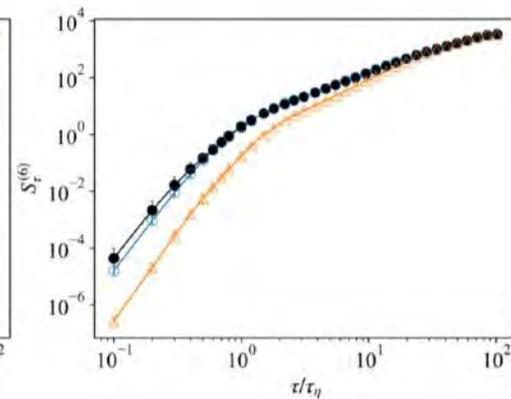
2nd order



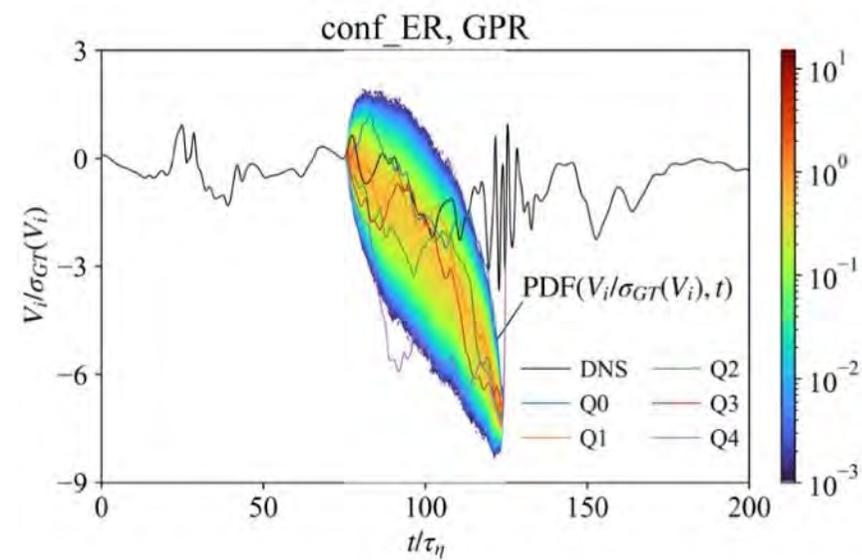
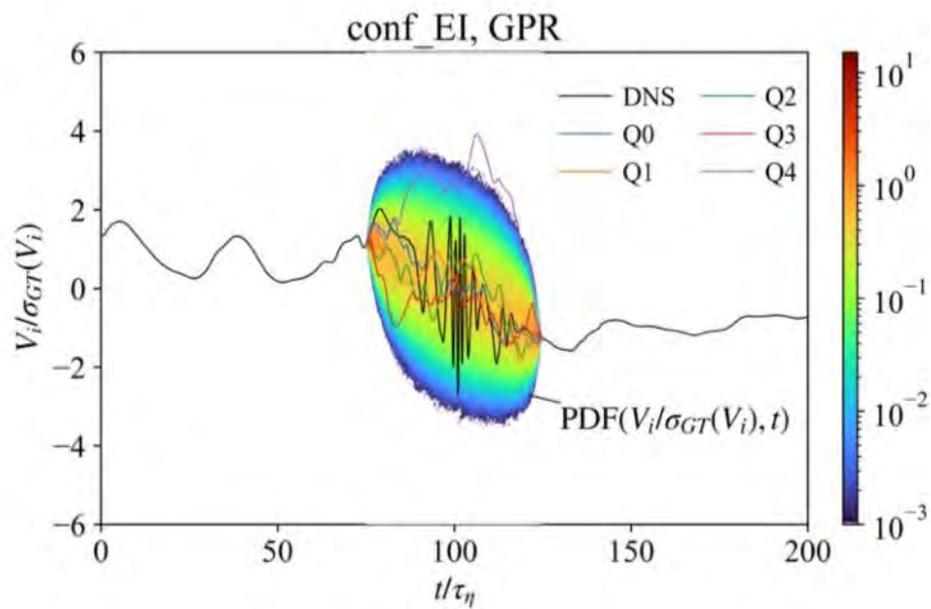
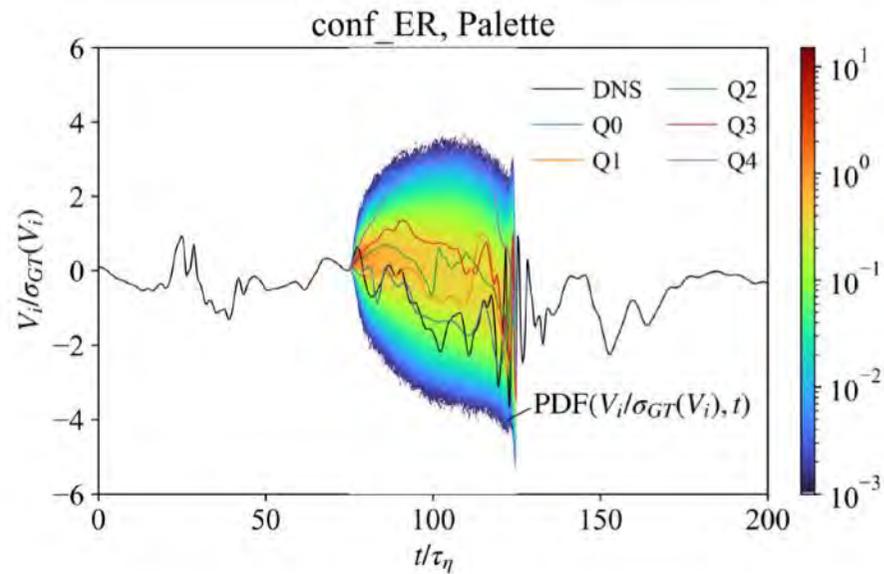
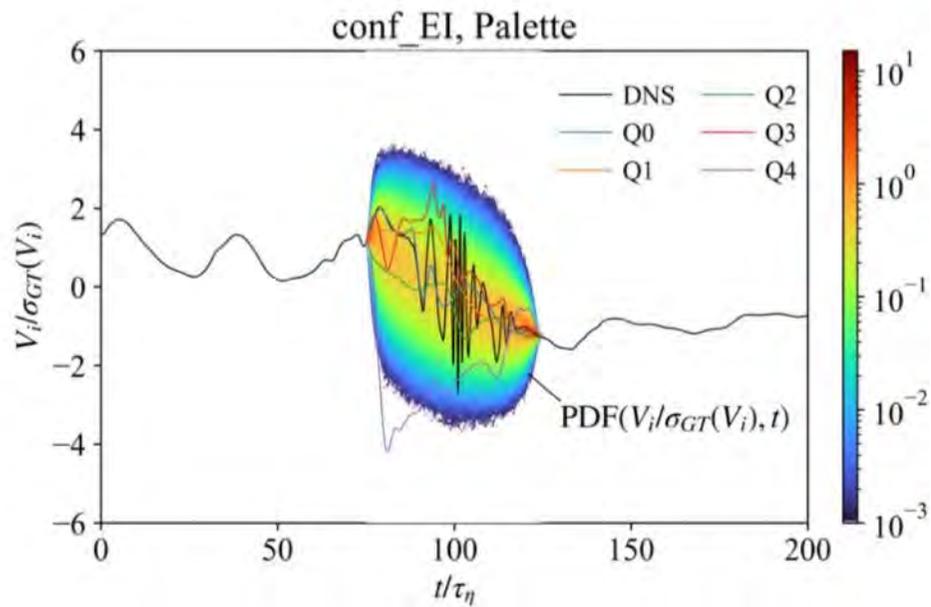
4th order



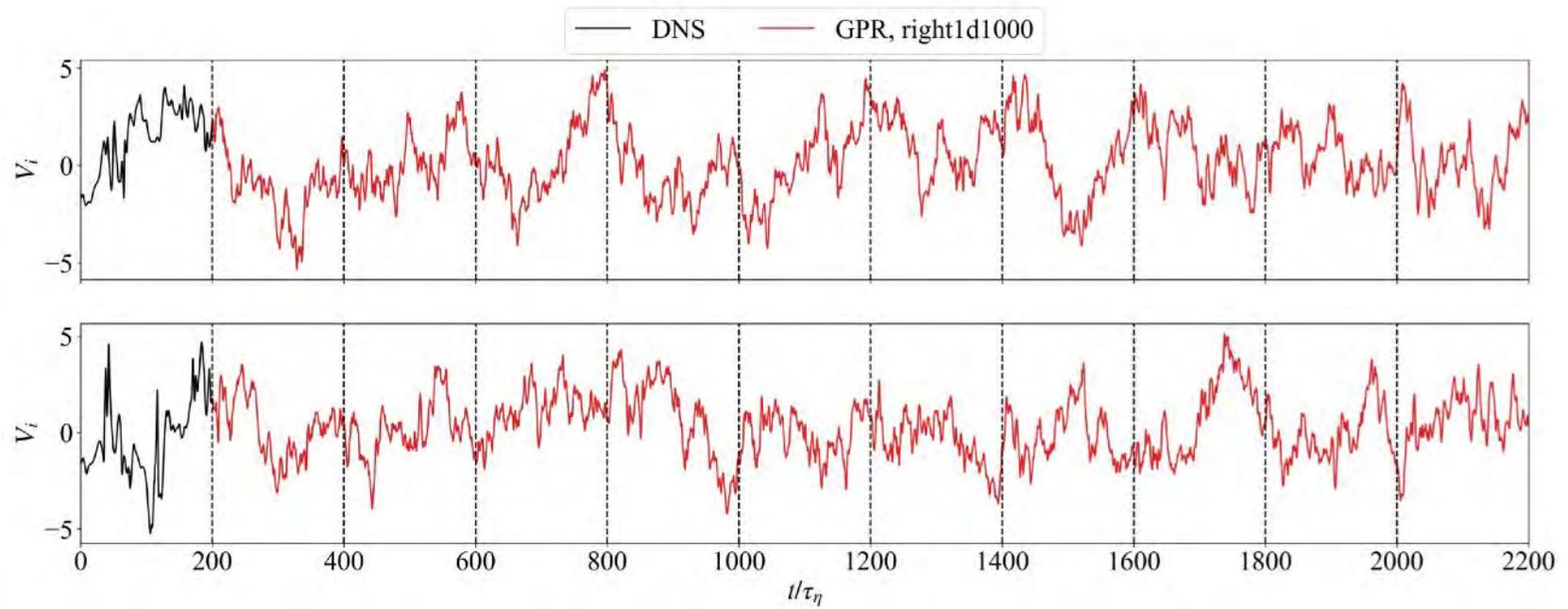
6th order

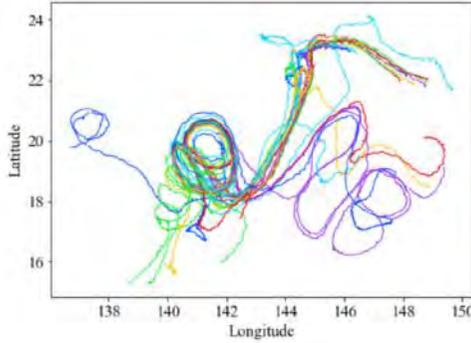
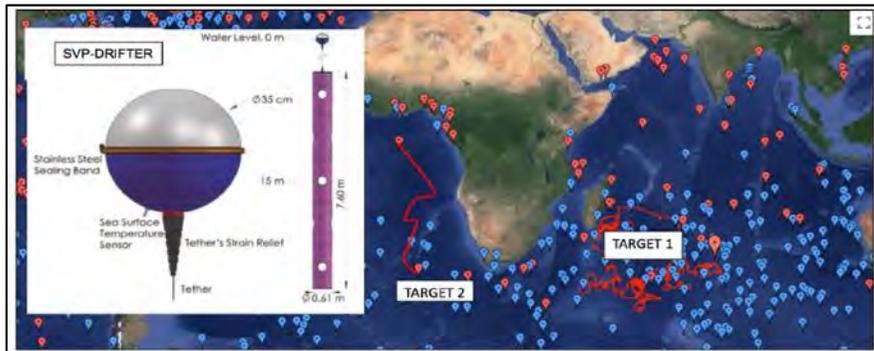


DNS
 Palette
 GPR



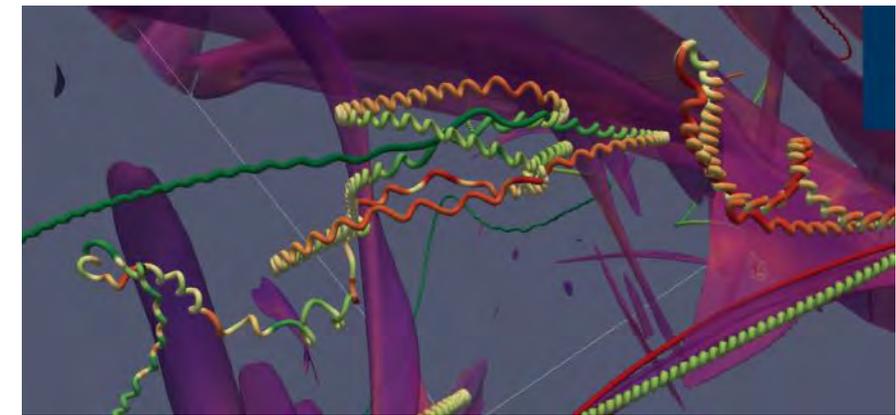
GENERATION IN THE LOOP





DRIFTERS FROM THE GLOBAL DRIFTER MAP
 SPATIAL TARGETS: (1) KEEP THE PROBES INSIDE A
 REGION AMONG TWO END-POINTS (ZERMELO PROBLEM).
 (2) KEEP THEM AT 15M DEPTH.

CHARGED-PARTICLES IN
 MHD
 with RL. Centurioni
 (UCSD, USA)



CHARGED-PARTICLES IN MHD
 with R. Grauer & J. Lubcke (Bochum U., GER)

**-WHAT-IF QUESTIONS: EXPLICABILITY OF THE GENERATED DATA, FEATURES RANKINGS,
 PHYSICS DISCOVERY**



Wavelet Score-Based Generative Modeling

<p>Florentin Guth Computer Science Department, ENS, CNRS, PSL University</p>	<p>Simon Coste Computer Science Department, ENS, CNRS, PSL University</p>
<p>Valentin De Bortoli Computer Science Department, ENS, CNRS, PSL University</p>	<p>Stéphane Mallat Collège de France, Paris, France Flatiron Institute, New York, USA</p>

[arxiv 2208.05003](https://arxiv.org/abs/2208.05003)